Maximal Net Baryon Density in the Energy Region Covered by NICA.

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	Equilibrium
π	$\exp\left[-rac{E_{\pi}}{T} ight]$
Ν	$\exp\left[-rac{E_N}{T}+rac{\mu_B}{T} ight]$
N	$\exp\left[-rac{E_N}{T}-rac{\mu_B}{T} ight]$
٨	$\exp\left[-\frac{E_{\Lambda}}{T}+\frac{\mu_{B}}{T}-\frac{\mu_{S}}{T} ight]$
$\overline{\Lambda}$	$\exp\left[-rac{E_{\Lambda}}{T}-rac{\mu_{B}}{T}+rac{\mu_{S}}{T} ight]$
К	$\exp\left[-rac{E_{\mathcal{K}}}{T}+rac{\mu_{\mathcal{S}}}{T} ight]$
ĸ	$\exp\left[-rac{E_{\kappa}}{T}-rac{\mu_{S}}{T} ight]$

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SPS data.

	Measurement		
Pb–Pb 158A GeV			
$(\pi^+ + \pi^-)/2.$	600±30		
K+	95 ±10		
K-	50 ± 5		
K ⁰ _S	60 ±12		
p	140±12		
p	10 ±1.7		
ϕ	7.6±1.1		
Ξ-	4.42±0.31		
Ξ-	0.74±0.04		
$\overline{\Lambda}/\Lambda$	0.2±0.04		



SPS data.

SPS: Freeze-Out Parameters:

 $T = 156.0 \pm 2.4 \text{MeV}$ $\mu_B = 239 \pm 12 \text{MeV}$

F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich Physical Review C64 (2001) 024901.



AGS data.

	Measurement		
Au–Au 11.6A GeV			
Participants	363±10		
K+	23.7±2.9		
K-	3.76±0.47		
π^+	133.7±9.9		
Λ	20.34±2.74		
p/π^+	1.234±0.126		
p	>0.0185±0.0018		



AGS data.

AGS: Freeze-Out Parameters:

 $T = 130.6 \pm 5.5 \text{MeV}$ $\mu_B = 594 \pm 26 \text{MeV}$

F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich Physical Review C64 (2001) 024901.



SIS data.

	Measurement			
Au–Au 1.7A GeV				
π^+/p	0.052±0.013			
K^+/π^+	0.003±0.00075			
π^-/π^+	2.05±0.51			
η/π^0	$0.018{\pm}0.007$			



SIS data.

SIS: Freeze-Out Parameters:

$$T = 49.7 \pm 1.1 \text{MeV}$$

 $\mu_B = 818 \pm 15 \text{MeV}$

J. C., H. Oeschler and K. Redlich) Physical Review C59, (1999) 1663.







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Chemical Freeze-Out: Criteria





Chemical Freeze-Out: Criteria





Chemical Freeze-Out: Criteria





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Chemical Freeze-Out





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Chemical Freeze-Out Temperature





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Chemical Freeze-Out μ_B





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 μ_B as a function of $\sqrt{s_{NN}}$

$$\mu_B(\sqrt{s}) = \frac{1.308 \text{ GeV}}{1 + 0.273 \text{ GeV}^{-1}\sqrt{s}}$$

This predicts at LHC $\mu_B \approx$ 1 MeV.

J. C., H. Oeschler, K. Redlich, S. Wheaton Phys. Rev. C73 034905 (2006)







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J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.



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Transition





Λ/π Ratio



THERMUS

S. Wheaton, J. Cleymans, M. Hauer

Comp. Phys. Comm. 180 (2009) 84-106



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Strangeness in Heavy Ion Collisions vs Strangeness in pp - collisions

Use the Wroblewski factor

$$\lambda_{m{s}} = rac{2\left< m{sar{m{s}}}
ight>}{\left< m{uar{m{u}}}
ight> + \left< m{dar{m{d}}}
ight>}$$

This is determined by the number of **newly** created quark – anti-quark pairs and **before** strong decays, i.e. before ρ 's and Δ 's decay.

Limiting values : $\lambda_s = 1$ all quark pairs are equally abundant, SU(3) symmetry. $\lambda_s = 0$ no strange quark pairs.



Maxima in particle ratios : K^+/π^+





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Maxima in particle ratios : K^+/π^+





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Maxima in particle ratios : K^+/π^+







R. Pisarski and L. McLerran



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J.C., H. Oeschler, K. Redlich, S. Wheaton, Phys. Lett. B615 (2005) 50-54

In the statistical model a rapid change is expected as the hadronic gas undergoes a transition from a baryon-dominated to a meson-dominated gas. The transition occurs at a temperature T = 151 MeV and baryon chemical potential $\mu_B = 327$ MeV corresponding to an incident energy of $\sqrt{s_{NN}} = 11$ GeV.



In conclusion, the roller-coaster seen in the particle ratios corresponds to a transition from a baryon-dominated to a meson-dominated hadronic gas. This transition occurs at a

- temperature T = 151 MeV,
- baryon chemical potential $\mu_B = 327$ MeV,
- energy $\sqrt{s_{NN}} = 11$ GeV.

In the statistical model this transition leads to peaks in the $\Lambda/\langle \pi \rangle$, K^+/π^+ , Ξ^-/π^+ and Ω^-/π^+ ratios.



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Thermal Model

The number of particles of type *i* is determined by:

$$E\frac{dN_i}{d^3p} = \frac{g_i}{(2\pi)^3} \int d\sigma_\mu p^\mu \exp\left(-\frac{p^\mu u_\mu}{T} + \frac{\mu_i}{T}\right)$$

Integrating this over all momenta

$$egin{aligned} N_i &= rac{g_i}{(2\pi)^3} \int d\sigma_\mu \int rac{d^3 p}{E} p^\mu \exp\left(-rac{p^\mu u_\mu}{T} + rac{\mu_i}{T}
ight) \ N_i &= \int d\sigma_\mu u^\mu n_i(T,\mu) \end{aligned}$$

or

If the temperature and chemical potential are unique along the freeze-out curve

$$N_i = n_i(T,\mu) \int d\sigma_\mu u^\mu$$

i.e. integrated (4π) multiplicities are the same as for a single fireball at rest (apart from the volume).

