

Precise predictions of low-energy QCD and their check by the DIRAC experiment

**99th Session of the JINR Scientific Council
19–20 January 2006**



Outline

- ➡ High-energy and low-energy QCD
- ➡ Precise predictions of low-energy QCD
- ➡ Experimental check of low-energy QCD predictions
- ➡ First lifetime measurement of the $\pi^+\pi^-$ -atom
- ➡ The new experiment on the investigation of $\pi^+\pi^-$ -atom and observation of πK -atoms at PS CERN
- ➡ Potentials of the DIRAC setup at SPS CERN, GSI CERN and J-PARC

DIRAC collaboration

75 Physicists from 18 Institutes



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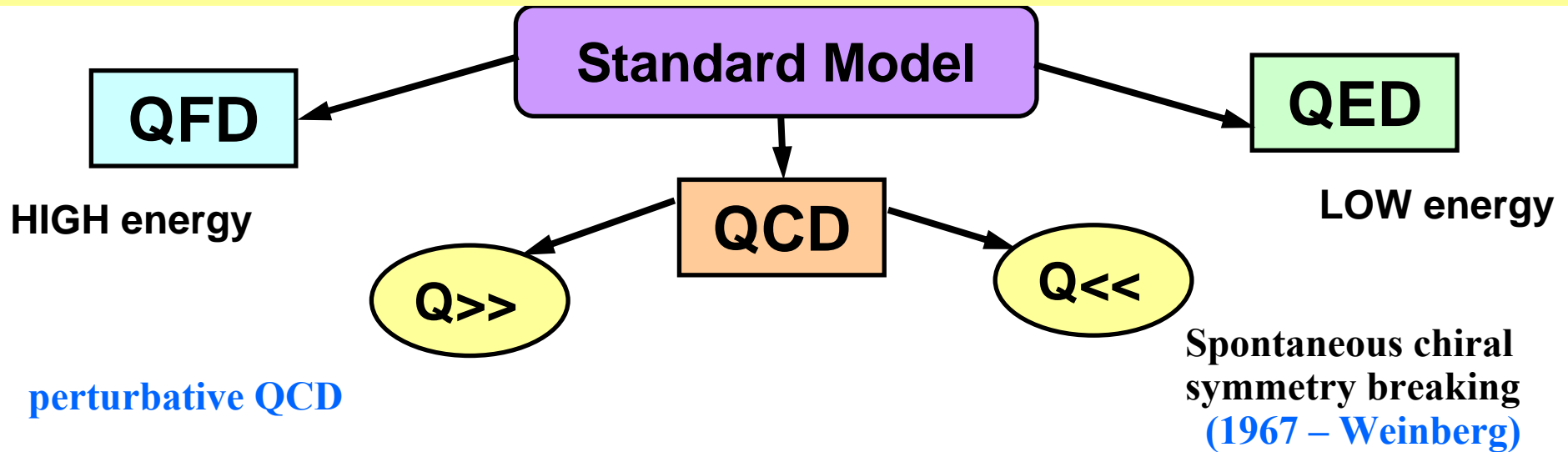
Bern University
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Theoretical motivation



QCD Lagrangian in presence of quark masses:

$$\mathcal{L}_{\text{QCD}(q,g)} = \mathcal{L}_{\text{sym}} + \mathcal{L}_{\text{break-sym}}$$

- high energy (small distance)
- “weak” interaction (asymptotic freedom)
- expansion in coupling

- low energy (large distance)
- strong interaction (confinement)
- expansion in momentum & mass

$$\mathcal{L}_{\text{eff}}(\pi, K, \eta) = \mathcal{L}_{\text{sym}} + \mathcal{L}_{\text{break-sym}}$$

M, for large Q, depends only on: \mathcal{L}_{sym}

M, for small Q, depends on both:
 \mathcal{L}_{sym} and $\mathcal{L}_{\text{break-sym}}$ and q-condensate

At low energies, QCD is replaced by an effective quantum field theory (ChPT)

formulated in terms of asymptotically observable fields like π , K, η

1979 – Weinberg

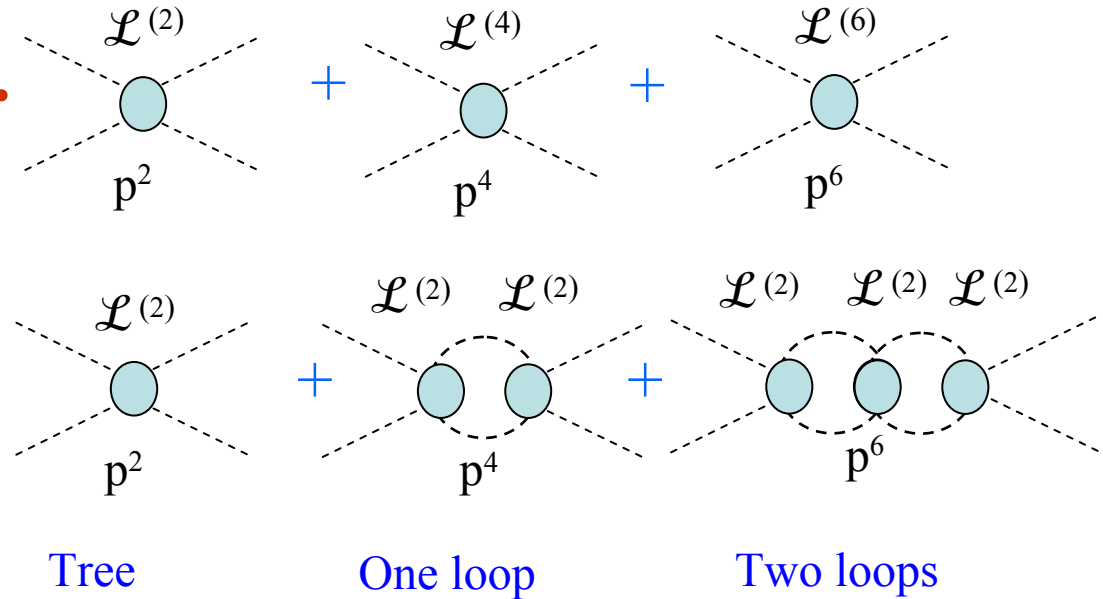
1984 – Gasser & Leutwyler

ChPT predictions for

$$\Delta = \left| a_0^0 - a_0^2 \right|$$

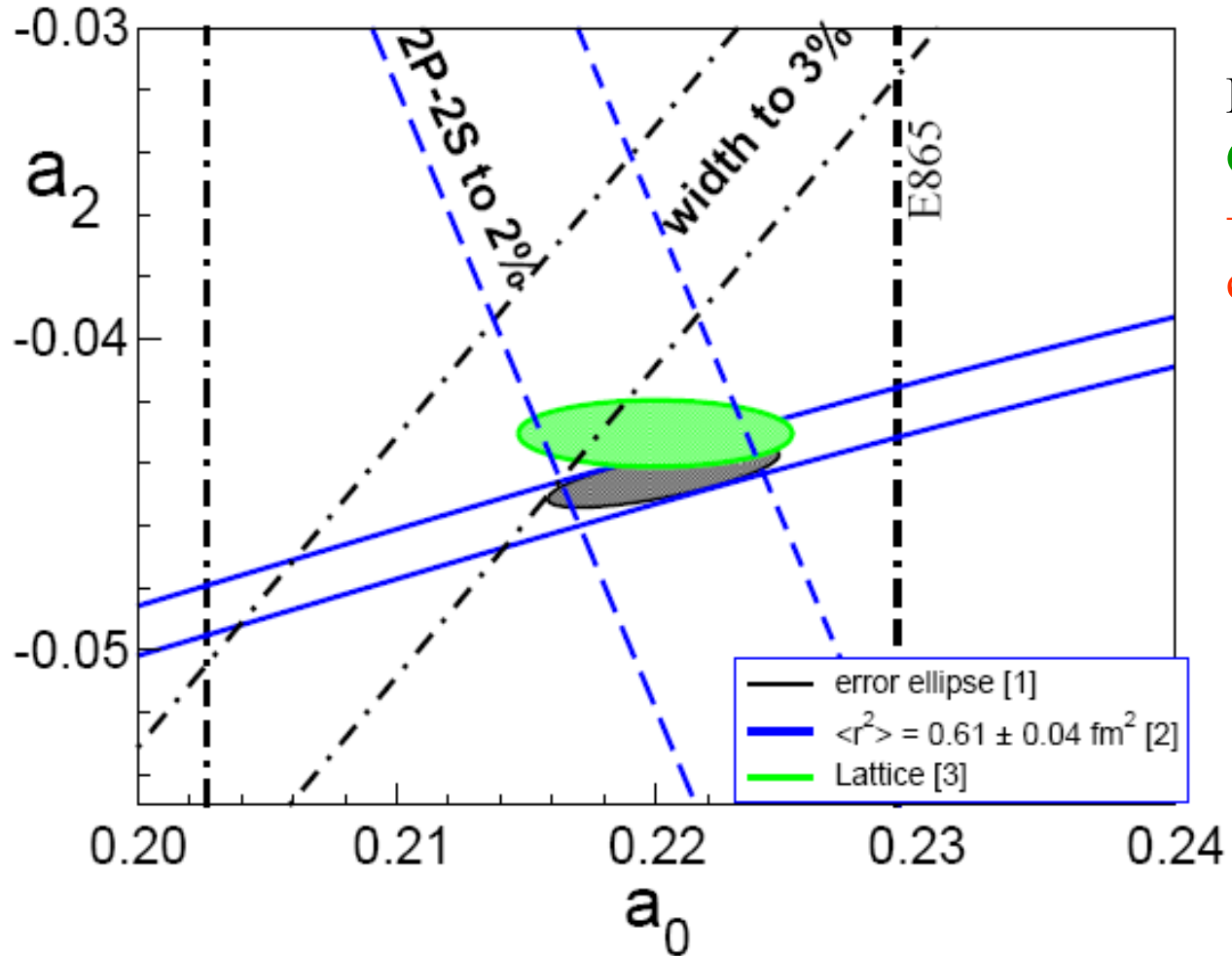
In ChPT effective Lagrangian \mathcal{L}_{eff} is constructed as an expansion in powers of external momenta and of quark masses

$$\mathcal{L}_{\text{eff}} = \mathcal{L}^{(2)} + \mathcal{L}^{(4)} + \mathcal{L}^{(6)} + \dots$$



1966 Weinberg (tree):	$\mathcal{L}^{(2)}$	$a_0 - a_2 = 0.20$	$a_0 = 0.159$	$a_2 = -0.045$
1984 Gasser-Leutwyler (1-loop):	$\mathcal{L}^{(4)}$	$a_0 - a_2 = 0.25 \pm 0.01$	$a_0 = 0.203$	$a_2 = -0.043$
1995 Knecht <i>et al.</i> (2-loop):	$\mathcal{L}^{(6)}$	Generalized ChPT		
1996 Bijnens <i>et al.</i> (2-loop):	$\mathcal{L}^{(6)}$	$a_0 - a_2 = 0.258 \pm (<3\%)$	$a_0 = 0.217$	$a_2 = -0.042$
2001 Colangelo <i>et al.</i> (& Roy):	$\mathcal{L}^{(6)}$	$a_0 - a_2 = 0.265 \pm 0.004 (1.5\%)$	$a_0 = 0.220$	$a_2 = -0.044$

Pions: $SU(2) \times SU(2)$



Dark ellipse \rightarrow Roy eqs. + ChPT
Green ellipse \rightarrow Roy eqs. + ChPT
+ some constants from Lattice
calculation (MILC Collaboration)

[1] Colangelo, Gasser and Leutwyler, 2001.

[2] Truong and Willey, 1989; Moussallam, 1999; Donoghue, Gasser and Leutwyler, 1990; Ynduráin 2003; Ananthanarayan *et al.* 2004.

[3] MILC Collaboration, 2004.

$\pi\pi$ scattering and quark condensate

In standard ChPT, quark condensate is LARGE:

$$\langle 0 | \bar{q}q | 0 \rangle \approx (-250 \text{ MeV})^3, \quad B = \lim_{\hat{m} \rightarrow 0} \left| \langle 0 | \bar{q}q | 0 \rangle \right| / F_\pi^2 \quad q = u, d$$

$$\longrightarrow m_\pi^2 = 2B\hat{m} + O(\hat{m}) \quad \text{with } B \text{ large, } \hat{m} \doteq (m_u + m_d) / 2$$

The linear term provides the dominant contribution to the π mass expansion:

$$(a_0 - a_2)_{\text{exp}} = (a_0 - a_2)_{\text{th}}$$

$$\text{Quark mass ratio} \rightarrow r = \frac{m_s}{\hat{m}} \approx 25$$

$$r = 25.7 \pm 2.6 \quad \text{Gasser and Leutwyler, 1985}$$

In Generalized ChPT [1, 2] quark condensate can be SMALL

$$\langle 0 | \bar{q}q | 0 \rangle \approx (-90 \text{ MeV})^3, \quad q = u, d$$

1. N.H. Fuchs, H. Sazdjian and J. Stern (1991)

2. M. Knecht *et al.* (1995)

$$\longrightarrow m_\pi^2 = 2B\hat{m} + 4A\hat{m}^2 + \dots \quad \text{where } B \text{ and } A \text{ are terms of the same order}$$

$$(a_0 - a_2)_{\text{exp}} > (a_0 - a_2)_{\text{th}}$$

$$r = \frac{m_s}{\hat{m}} < 25$$

$$(a_0 - a_2)_{\text{exp}} < (a_0 - a_2)_{\text{th}}$$

No explanation exist

$\pi\pi$ scattering lengths

Present low energy QCD predictions:

$$a_0 = 0.220 \pm 0.005 (2.3\%)$$

$$a_2 = -0.0444 \pm 0.0010 (2.3\%)$$

$$a_0 - a_2 = 0.265 \pm 0.004 (1.5\%)$$

First result:

L. Rosselet *et al.*,
Phys. Rev. D15 (1977) 574

$$a_0 = 0.28 \pm 0.05 (18\%) \text{ using Roy eqs.}$$

Results from E865/BNL: $K \rightarrow \pi^+\pi^-e^+\nu_e (K_{e4})$

S.Pislak *et al.*, Phys. Rev. Lett. 87 (2001) 221801
using Roy eqs.

$$a_0 = 0.203 \pm 0.033 (16\%)$$

$$a_2 = -0.055 \pm 0.023 (42\%)$$

using Roy eqs. and ChPT constraints $a_2 = f_{ChPT}(a_0)$

$$a_0 = 0.216 \pm 0.013 (stat) \pm 0.004(syst) \pm 0.002 (theor)$$

$$\delta a_0 = \pm 6\% (stat) \pm 2\%(syst) \pm 1\% (theor)$$

Results from NA48/2: $K^+ \rightarrow \pi^0\pi^0\pi^+$

$$(a_0 - a_2)m_\pi = 0.268 \pm 0.010(stat) \pm 0.004(syst)$$

$$\delta(a_0 - a_2) = \pm 3.7\%(stat) \pm 1.5\%(syst) \pm 5\%(theor)$$

$$(a_0 - a_2)m_\pi = 0.264 \pm 7.5\%(stat) \begin{matrix} +3\% \\ -8\% \end{matrix} (syst)$$

$$\delta(a_0 - a_2) = \pm 5\%(stat) \begin{matrix} +3\% \\ -8\% \end{matrix} (syst)$$

DIRAC current results, 2001 data

DIRAC expected results, 2001–2003 data

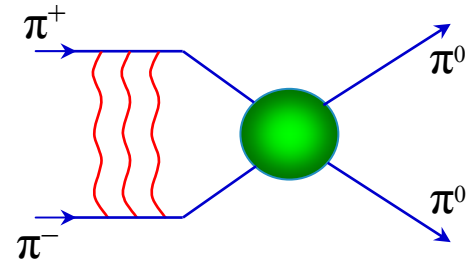
Upgraded DIRAC

$$\delta(a_0 - a_2) = \pm 2\%(stat) \pm 1\%(syst) \pm 1\%(theor)$$

Theoretical limits

1. $A_{2\pi}$ time of life

$$A_{2\pi} \rightarrow \pi^0 \pi^0 \quad \Gamma(\pi^0 \pi^0) = R_\pi (a_0 - a_2)^2 (1 + \delta_\pi)$$



H. Jalloul, H.Sazdjian 1998

M.A. Ivanov et al. 1998

A. Gashi et al. 2002

J. Gasser et al. 2001

$$\rightarrow \delta_\pi = (5.8 \pm 1.2) \cdot 10^{-2}$$

Current limit for accuracy in scattering lengths measurement from the $A_{2\pi}$ lifetime

$$\frac{\Delta |a_0 - a_2|}{|a_0 - a_2|} = 0.6\%$$

2. $A_{2\pi}$ interaction with matter

L.Afanasyev, G.Baur, T.Heim, K.Hencken, Z.Halabuka, A.Kotsinyan, S.Mrowczynski, C.Santamarina, M.Schumann, A.Tarasov, D.Trautmann, O.Voskresenskaya from Basel, JINR and CERN

Current limit for accuracy in scattering lengths measurements due to accuracy in $P_{br}(\tau)$

$$\frac{\Delta |a_0 - a_2|}{|a_0 - a_2|} = 1.2\%$$

This value will be reduced by a factor of 2

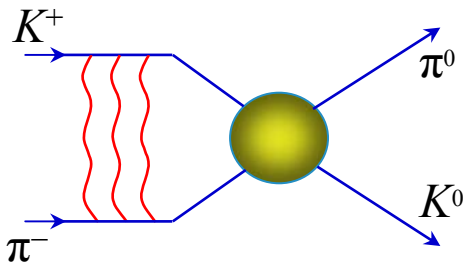
Theoretical limits

$\Lambda_{K^+\pi^-}$ and $\Lambda_{K^+\pi^-}$ time of life

$$A_{K^+\pi^-} \rightarrow \pi^0 K^0$$

$$\Gamma(\pi K) = R_K |a_{1/2} - a_{3/2}|^2 (1 + \delta_K)$$

$$A_{\pi^+K^-} \rightarrow \pi^0 \bar{K}^0$$



$$\delta_K = (4.0 \pm 2.2) \cdot 10^{-2} \frac{\Delta |a_{1/2} - a_{3/2}|}{|a_{1/2} - a_{3/2}|} = 1.1\%$$

J. Schweizer (2004)

Production of ponium

Atoms are Coulomb bound state of two pions produced in one proton-nucleus collision

$$\frac{d\sigma_{nlm}^A}{d\vec{P}} = (2\pi)^3 \frac{E_A}{M_A} \left| \psi_{nlm}^{(C)}(0) \right|^2 \frac{d\sigma_s^0}{d\vec{p}_+ d\vec{p}_-} \Big|_{\vec{p}_+ = \vec{p}_-}$$

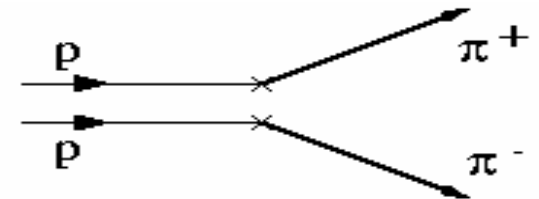
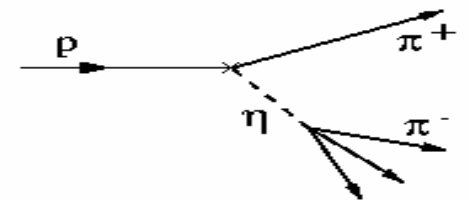
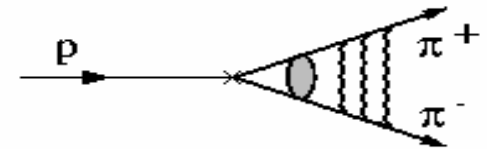
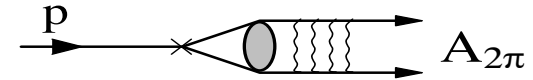
Background processes:

Coulomb pairs. They are produced in one proton nucleus collision from fragmentation or short lived resonances and exhibit Coulomb interaction in the final state

$$\frac{d^2\sigma_C}{d\vec{p}_+ d\vec{p}_-} = A_C(q) \frac{d\sigma_s^0}{d\vec{p}_+ d\vec{p}_-}, \quad A_C(q) = \frac{2\pi m_\pi \alpha / q}{1 - \exp(-2\pi m_\pi \alpha / q)}$$

Non-Coulomb pairs. They are produced in one proton nucleus collision. At least one pion originates from a long lived resonance. No Coulomb interaction in the final state

Accidental pairs. They are produced in two independent proton nucleus collision. They do not exhibit Coulomb interaction in the final state



Method of $A_{2\pi}$ observation and lifetime measurement

L. Nemenov, Sov. J. Nucl. Phys. (1985)

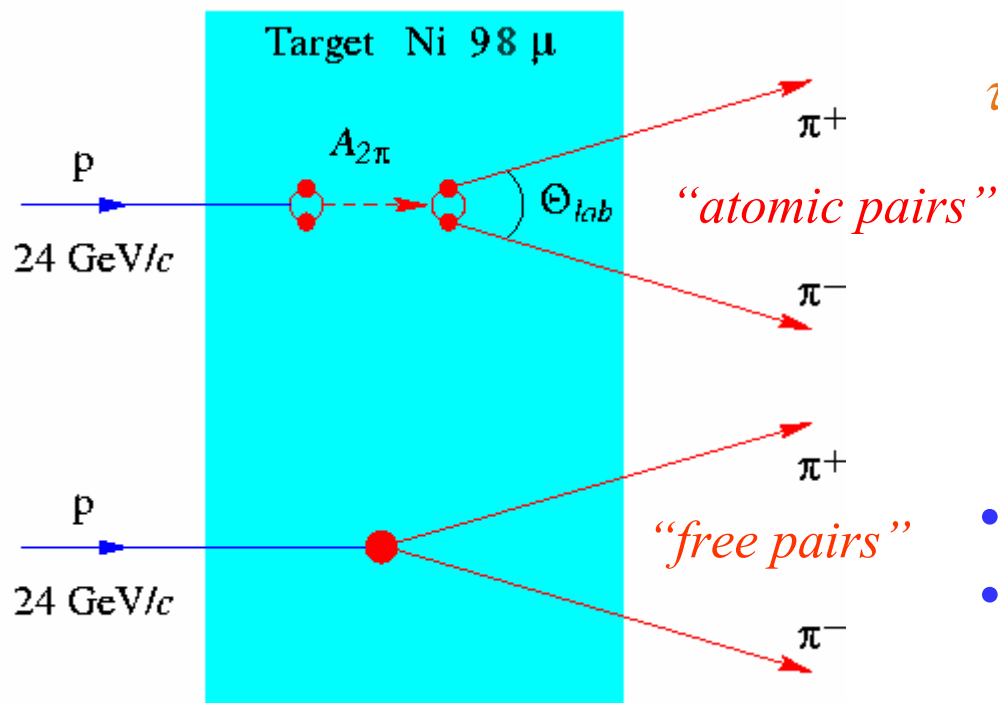
$\tau(A_{2\pi})$ too small to be measured directly

e.m. interaction of $A_{2\pi}$ in the target

$$A_{2\pi} \rightarrow \pi^+ \pi^-$$

$$Q < 3 \text{ MeV}/c \quad E_+ \approx E_- \quad \Theta_{lab} < 2.5 \text{ mrad}$$

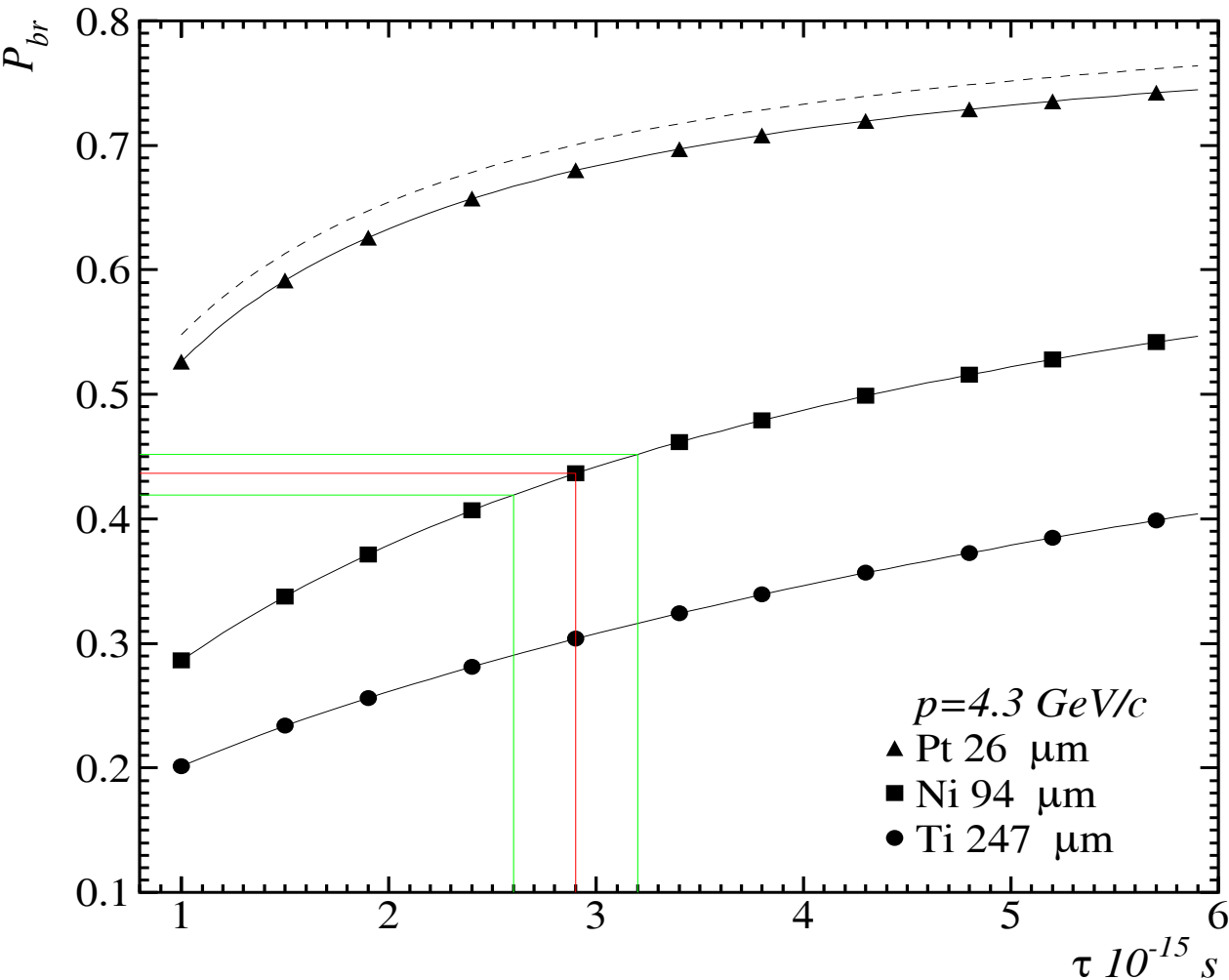
- *Coulomb from short-lived sources*
- *non-Coulomb from long-lived sources*



First observation of $A_{2\pi}$ have been done by the group from JINR, SINP MSU and IHEP at U-70 Protvino Afanasyev L.G. et al., Phys.Lett.B, 1993.

Break-up probability

Solution of the transport equations provides one-to-one dependence of the measured break-up probability (P_{br}) on pionium lifetime τ



All targets have the same thickness in radiation lengths $6.7 \cdot 10^{-3} X_0$

There is an optimal target material for a given lifetime

The detailed knowledge of the cross sections (Afanasyev&Tarasov; Trautmann et al) (Born and Glauber approach) together with the accurate description of atom interaction dynamics (including density matrix formalism) permits us to know the curves within 1%.

Energy splitting

Annihilation: $A_{2\pi} \rightarrow \pi^0 \pi^0$ $1/\tau = W_{\text{ann}} \sim (a_0 - a_2)^2$

Energy Splitting between np - ns states in $A_{2\pi}$ atom

$$\Delta E_n \equiv E_{ns} - E_{np}$$

$$\Delta E_n \approx \Delta E_n^{\text{vac}} + \Delta E_n^s \quad \Delta E_n^s \sim 2a_0 + a_2$$

For $n=2$ $\Delta E_2^{\text{vac}} = -0.107 \text{ eV}$ *from QED calculations*

$\Delta E_2^s \approx -0.45 \text{ eV}$ *numerical estimated value from ChPT*

$$a_0 = 0.220 \pm 0.005$$

$$a_2 = -0.0444 \pm 0.0010$$

(2001) G. Colangelo, J. Gasser and H. Leutwyler

$$\Rightarrow \Delta E_2 \approx -0.56 \text{ eV}$$

(1979) A. Karimkhodzhaev and R. Faustov

(1983) G. Austen and J. de Swart

(1986) G. Efimov *et al.*

(1999) A. Gashi *et al.*

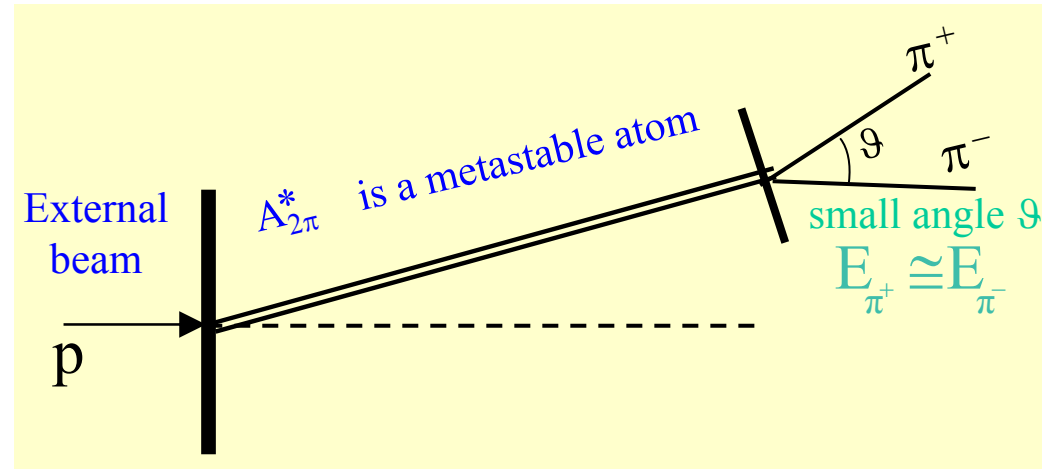
(2000) D. Eiras and J. Soto

Measurement of τ and ΔE allows one to obtain a_0 and a_2 separately

Metastable atoms

For $p_A = 5.6 \text{ GeV}/c$ and $\gamma = 20$

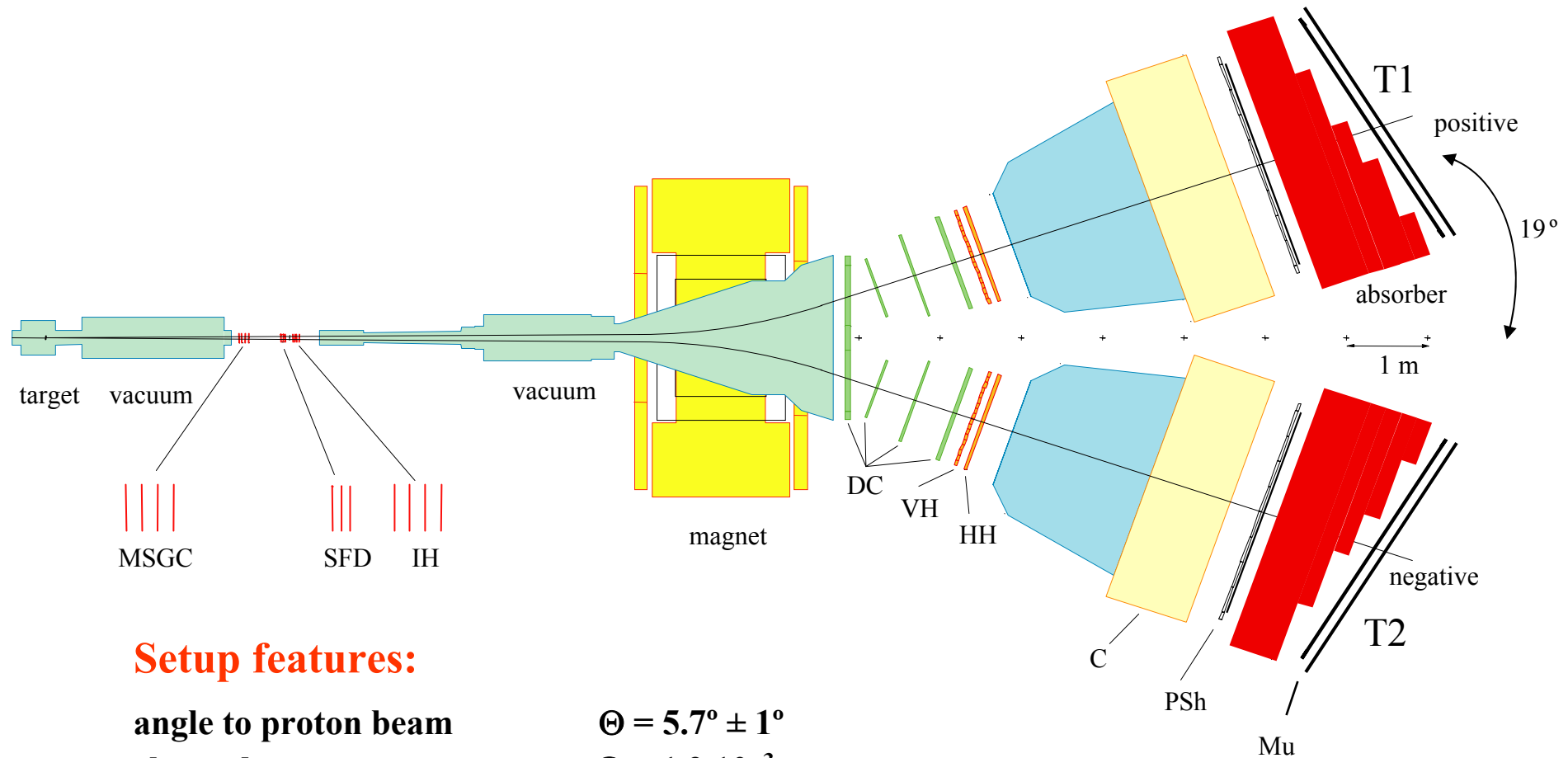
$$\left\{ \begin{array}{ll} \tau_{1s} = 2.9 \times 10^{-15} \text{ s}, & \lambda_{1s} = 1.7 \times 10^{-3} \text{ cm} \\ \tau_{2s} = 2.3 \times 10^{-14} \text{ s}, & \lambda_{2s} = 1.4 \times 10^{-2} \text{ cm} \\ \tau_{2p} = 1.17 \times 10^{-11} \text{ s}, & \lambda_{2p} = 7 \text{ cm} \\ & \lambda_{3p} \approx 23 \text{ cm} \\ & \lambda_{4p} \approx 54 \text{ cm} \end{array} \right.$$



Probabilities of the $A_{2\pi}$ breakup (Br) and yields of the long-lived states for different targets provided the maximum yield of summed population of the long-lived states: $\Sigma(l \geq 1)$

Target Z	Thickness Mm	Br	Σ ($l \geq 1$)	$2p_0$	$3p_0$	$4p_0$	Σ ($l=1, m=0$)
04	100	4.45%	5.86%	1.05%	0.46%	0.15%	1.90%
06	50	5.00%	6.92%	1.46%	0.51%	0.16%	2.52%
13	20	5.28%	7.84%	1.75%	0.57%	0.18%	2.63%
28	5	9.42%	9.69%	2.40%	0.58%	0.18%	3.29%
78	2	18.8%	10.5%	2.70%	0.54%	0.16%	3.53%

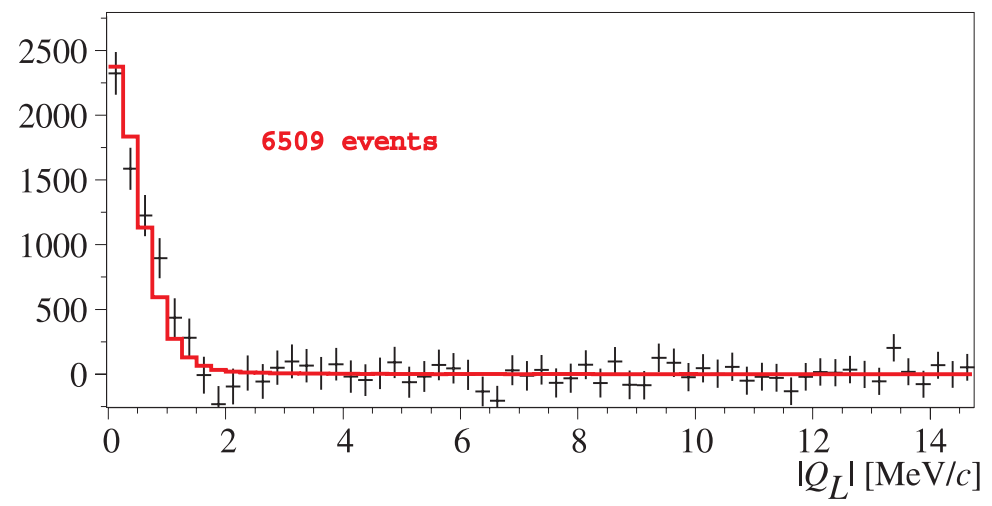
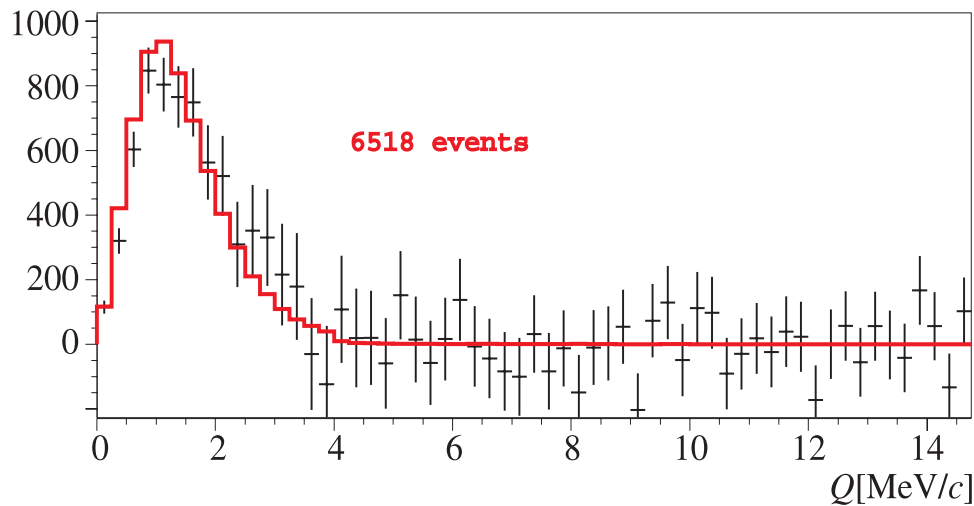
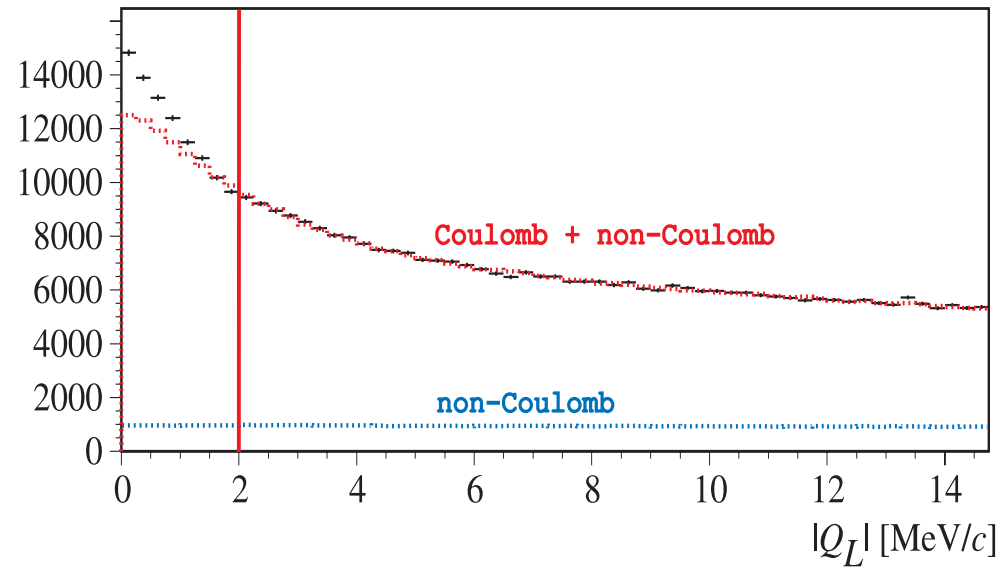
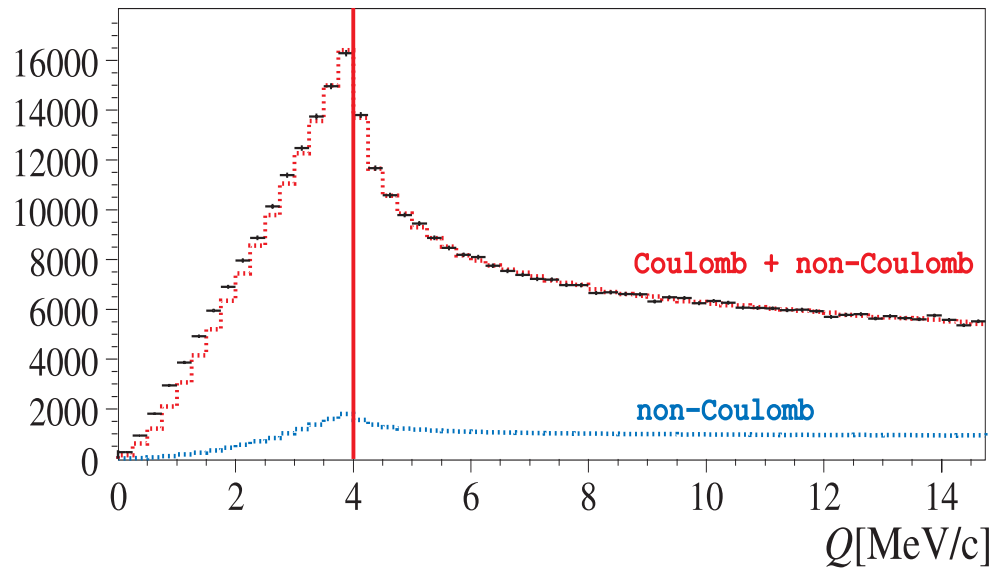
DIRAC set-up



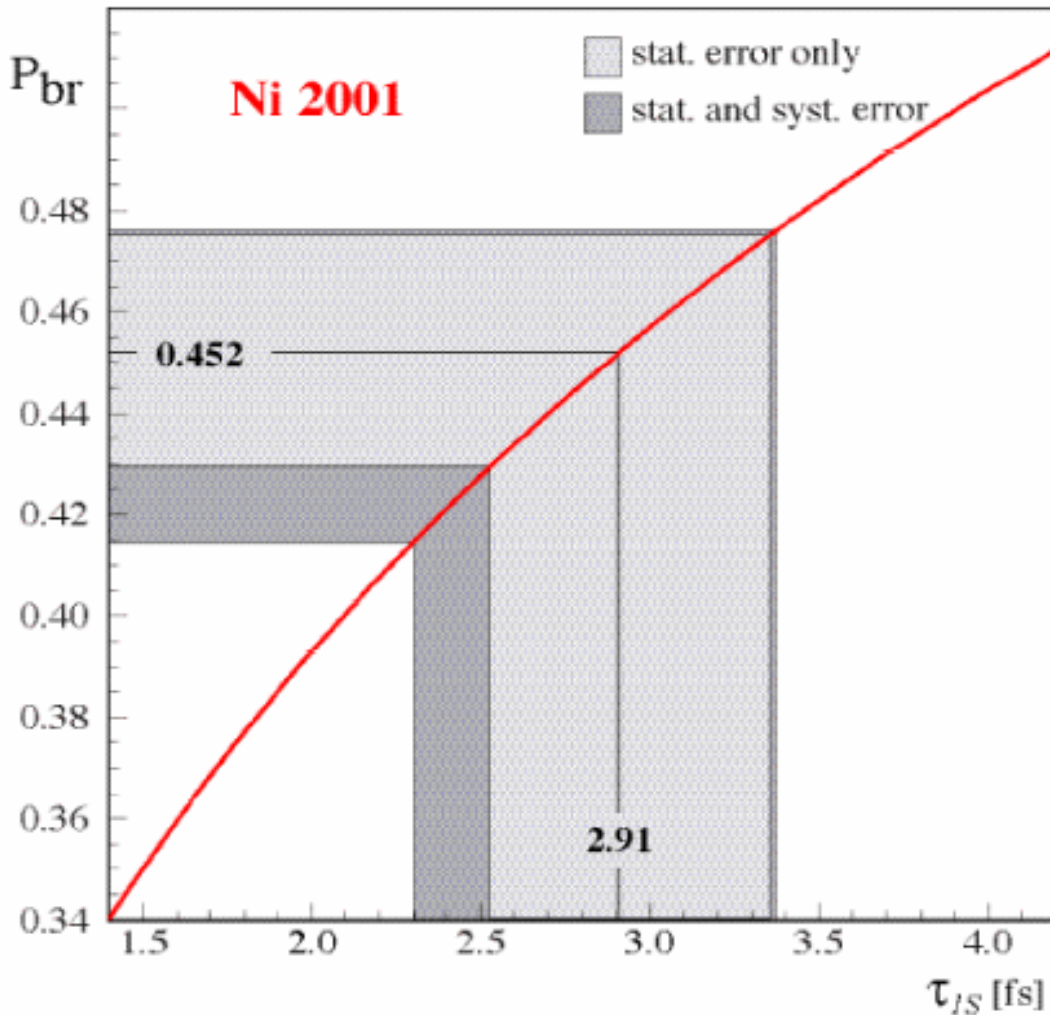
Setup features:

angle to proton beam	$\Theta = 5.7^\circ \pm 1^\circ$
channel aperture	$\Omega = 1.2 \cdot 10^{-3} \text{sr}$
dipole magnet	$B = 1.65 \text{ T}, BL = 2.2 \text{ Tm}$
momentum range	$1.2 \leq p_\pi \leq 8 \text{ GeV}/c$
momentum resolution	$\Delta p/p \approx 3 \cdot 10^{-3}$
resolution on relative momentum	$\sigma Q_x \approx \sigma Q_y \leq 0.5 \text{ MeV}/c,$ and $\sigma Q_L \approx 0.5 \text{ MeV}/c$

Atomic pairs



Lifetime of Pionium



Result from DIRAC:

$$\tau = \left(2.91^{+0.45}_{-0.38} \right)_{stat} \left({}^{+0.19}_{-0.49} \right)_{syst} \times 10^{-15} \text{ s}$$

$$|a_0 - a_2| = 0.264^{+0.033}_{-0.020} m_{\pi}^{-1}$$

ChPT prediction:

$$\tau = (2.9 \pm 0.1) \times 10^{-15} \text{ s}$$

$$a_0 - a_2 = 0.265 \pm 0.004$$

DIRAC analysis

Results for the lifetime:

$$\tau_{1S} = 2.91 \left. \begin{array}{l} +0.45 \\ -0.38 \end{array} \right\}_{stat} \left. \begin{array}{l} +0.19 \\ -0.49 \end{array} \right\}_{syst} = 2.91 \begin{array}{l} +0.49 \\ -0.62 \end{array} [fs]$$

$$\tau_{1S}^{ChPT} = 2.9 \pm 0.1 [fs]$$

Result for scattering lengths:

$$|a_0 - a_2| = 0.264 \begin{array}{l} +0.033 \\ -0.020 \end{array} [m_\pi^{-1}]$$

$$|a_0 - a_2|_{ChPT} = 0.265 \pm 0.004 [m_\pi^{-1}]$$

Improvements with full statistics

Number of Atomic pairs (approx.)

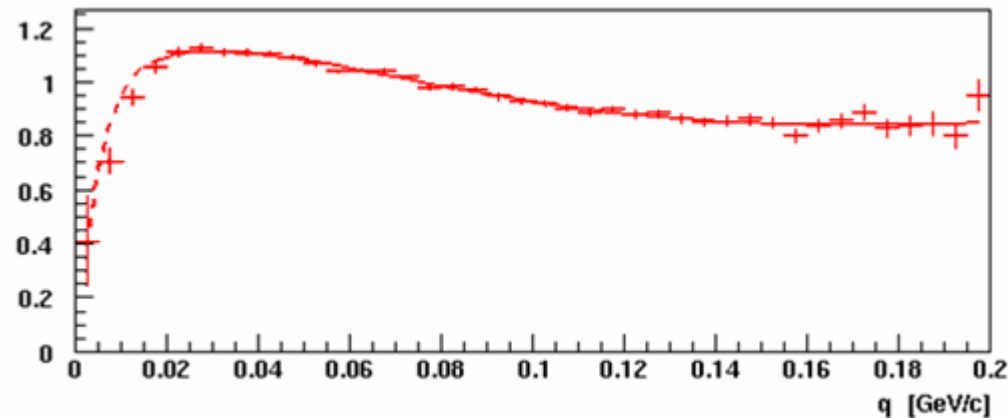
	Pt1999 24 GeV	Ni2000 24 GeV	Ti2000 24 GeV	Ti2001 24 GeV	Ni2001 24 GeV	Ni2002 20 GeV	Ni2002 24 GeV	Ni2003 20 GeV	Sum
Sharp selection	280	1300	900	1500	6500	3000	4500	1400	19400
Downstream only									27000

$$\frac{\sigma_{P_{br}}}{P_{br}} \Big|_{stat}^{now} = 0.051 \Rightarrow \frac{\sigma_{P_{br}}}{P_{br}} \Big|_{stat}^{full\ statistics} = 0.03 \Rightarrow \frac{\delta |a_0 - a_2|}{a_0 - a_2} \Big|_{stat} = 5\%$$

as in the project

Finite-size effects

CF($\pi^-\pi^-$) arbitrary normalization



Simulation vs fit of DIRAC $\pi^-\pi^-$ CF

- simulation $N_\omega(\pi^-\pi^-) = 19.2\%$
 - fit result $N_\omega(\pi^-\pi^-) = 21 \pm 7\%$
- ⇒ good description of ω pairs by UrQMD

In $\pi^+\pi^-$ system finite-size effect induces shift in P_{br}

✓ UrQMD simulation

$$N_\omega(\pi^+\pi^-) = 15\% \Rightarrow \delta P_{br} \sim 2\% \Rightarrow \delta\tau \sim 5\%$$

✓ upper limit at 1σ of $\pi^-\pi^-$ fit

$$N_\omega(\pi^+\pi^-) = 20\% \Rightarrow \delta P_{br} \sim 3\% \Rightarrow \delta\tau \sim 7.5\%$$

⇒

Systematic shift in τ measurement from finite-size effect $< 10\%$

i.e. less than present DIRAC statistical error in τ .

Expected shift with multi-layer target in future DIRAC 5 times less

What new will be known if πK scattering length will be measured?

The measurement of s-wave πK scattering length would test our understanding of chiral $SU(3)_L \times SU(3)_R$ symmetry breaking of QCD (u, d and s), while the measurement of $\pi\pi$ scattering length checks only $SU(2)_L \times SU(2)_R$ symmetry breaking (u, d).

This is the main difference between $\pi\pi$ and πK scattering!

πK scattering

I. ChPT predicts s-wave scattering lengths:

$$a_0^{1/2} = 0.19 \pm 0.2 \quad a_0^{3/2} = -0.05 \pm 0.02$$

$\mathcal{L}^{(2)}$, $\mathcal{L}^{(4)}$ and 1-loop

V. Bernard, N. Kaiser,
U. Meissner. – 1991

$$a_0^{1/2} - a_0^{3/2} = 0.23 \pm 0.01$$

A. Rossel. – 1999

$\mathcal{L}^{(2)}$, $\mathcal{L}^{(4)}$, $\mathcal{L}^{(6)}$ and 2-loop

J. Bijnens, P. Talaver. – April 2004

II. Roy-Steiner equations:

$$a_0^{1/2} - a_0^{3/2} = 0.269 \pm 0.015$$

III. $A_{\pi K}$ lifetime:

$$A_{\pi^+ K^-} \rightarrow \pi^0 \bar{K}^0 \quad (A_{K^+ \pi^-} \rightarrow \pi^0 K^0)$$

$$\Gamma(\pi^0 \bar{K}^0) \sim |a_0^{1/2} - a_0^{3/2}|^2 \quad \text{precision} \sim 1\% \quad \text{J. Schweizer. – 2004}$$

$$\tau = (3.7 \pm 0.4) \cdot 10^{-15} \text{ s}$$

Main goals and time scale for the $A_{2\pi}$ and $A_{\pi K}$ experiments

Manufacture of all new detectors and electronics: 18 months
Installation of new detectors: 3 months

2006

Test of the Upgraded setup and calibration: 4 months
Observation $A_{2\pi}$ in the long-lived states.

2007 and 2008

Measurement of $A_{2\pi}$ lifetime: 12 months

In this time 86000 $\pi\pi$ atomic pairs will be collected to estimate $A_{2\pi}$ lifetime with precision of:

$$\frac{\sigma_{\tau}}{\tau} = 6\%, \quad \frac{\sigma(a_0 - a_2)}{a_0 - a_2} = 3\%$$

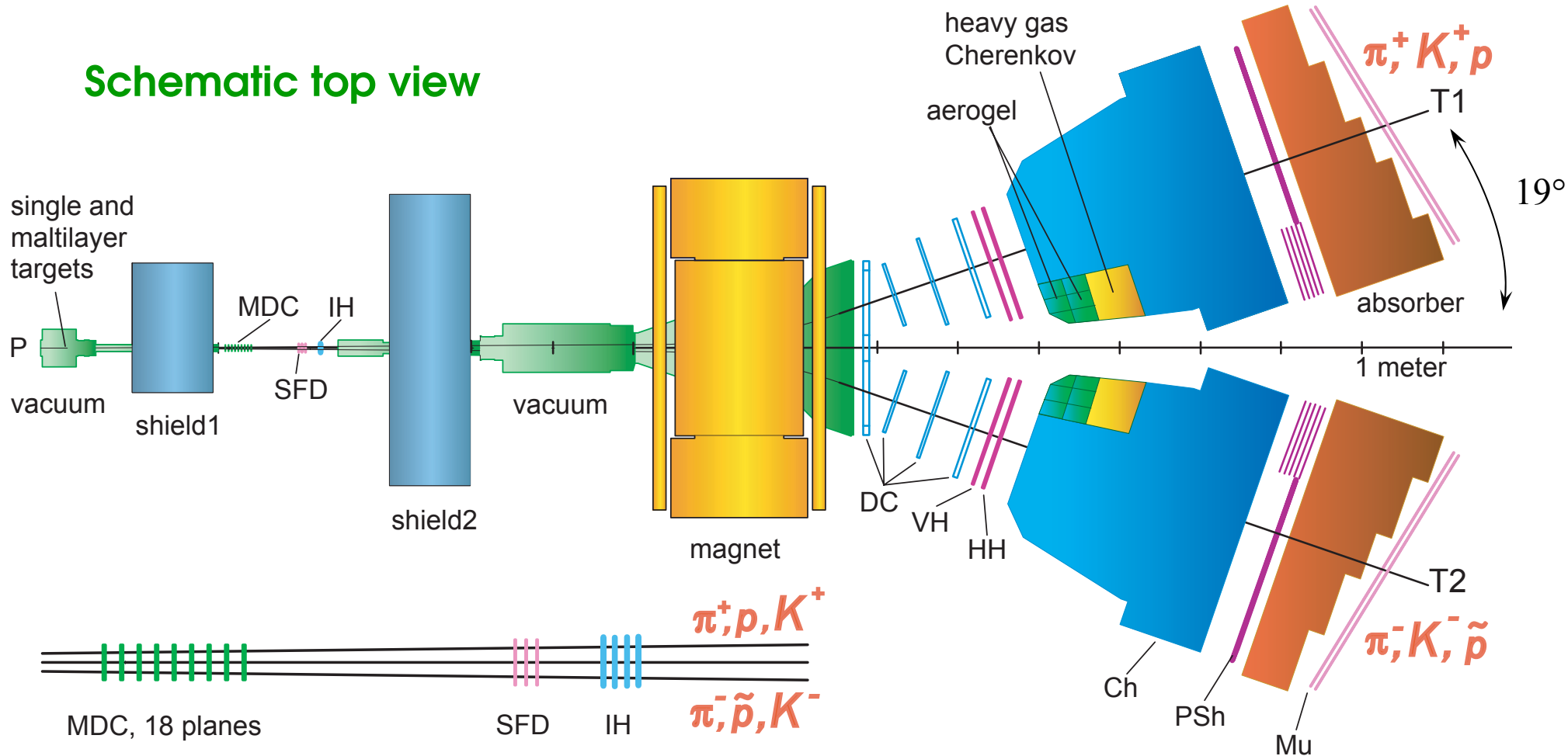
At the same time we also plan to observe $A_{\pi K}$ and $A_{K\pi}$; to detect 5000 πK atomic pairs to estimate $A_{\pi K}$ lifetime with precision of:

$$\frac{\sigma_{\tau}}{\tau} = 20\%, \quad \frac{\sigma(a_{1/2} - a_{3/2})}{a_{1/2} - a_{3/2}} = 10\%$$

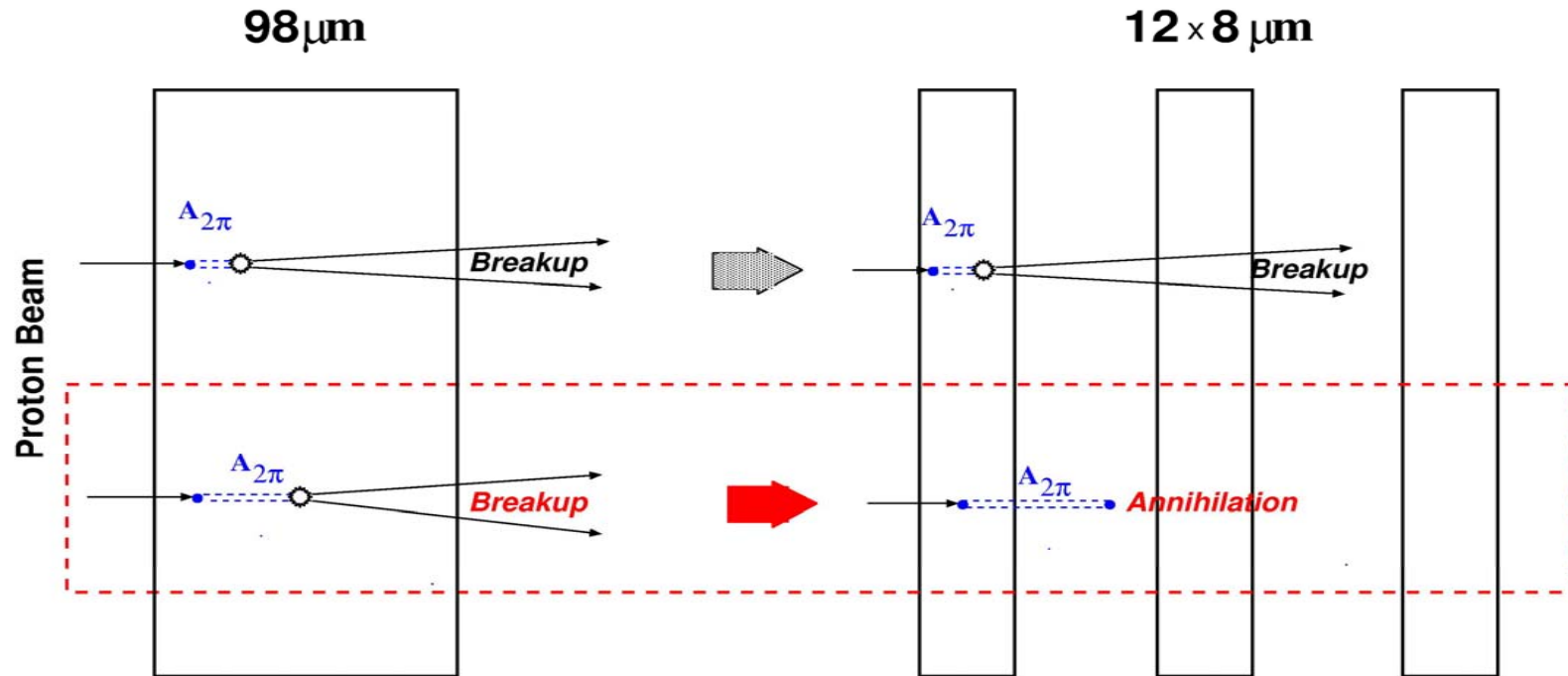
This estimation of the beam time is based on the $A_{2\pi}$ statistics collected in 2001 and on the assumption of having 2.5 spills per supercycle during 20 hours per day.

Upgrade DIRAC experimental set-up description

Schematic top view



Dual Target Method



- Single/Multilayer target comparison:
 - Same amount of multiple scattering
 - Same background (CC, NC, ACC)
 - Same number of produced $A_{2\pi}$, but lower number of dissociated pairs

Yields of atoms as a function of the proton beam momentum

Yields of pion pairs and atoms for the 24 GeV proton beam per one pNi-interaction at $\Theta_{\text{lab}}=5.7^\circ$

P, GeV/c	$\pi^+ \pi^-$	$A_{2\pi}$	$A_{2\pi} / \pi^+ \pi^-$	$A_{\pi K} + A_{K\pi}$	$(A_{\pi K} + A_{K\pi}) / \pi^+ \pi^-$
24	2.1×10^{-2}	0.95×10^{-9}	4.4×10^{-8}	0.83×10^{-10}	0.39×10^{-8}

Relative yields of pion pairs and atoms as a function of the proton beam momentum

	P, GeV/c	$\pi^+ \pi^-$	$A_{2\pi}$	$A_{2\pi} / \pi^+ \pi^-$	$A_{\pi K} + A_{K\pi}$	$(A_{\pi K} + A_{K\pi}) / \pi^+ \pi^-$	Duty factor
PS CERN	24	1	1	1	1	1	1(0.06)
GSI (SIS100)	30	1.2	1.4	1.14	1.5	1.26	8.4
J-PARC	50	1.6	2.2	1.43	2.8	1.74	3.3
GSI (SIS200)	60	1.8	2.6	1.52	3.5	1.91	8.4
GSI (SIS300)	90	2.0	3.4	1.72	4.6	2.30	8.4
SPS CERN	450	3.1	12	3.7	13.5	4.3	4.0

Expected accuracy for $\pi\pi$ -scattering

Estimation of error sources in $\Delta/a_0 - a_2 / |a_0 - a_2|$ based on data taken with the upgraded DIRAC setup during 12 months (20h/day)

Single-layer target

	Number of atomic pairs n_A	Relative statistical error $(a_0 - a_2)$	Relative theoretical error of $(a_0 - a_2)$ from the ratio $\tau = f(a_0 - a_2)$	Relative theoretical error of $(a_0 - a_2)$ from the ratio $P_{br} = \varphi(\tau)^{(*)}$	Error from non point-like production
PS CERN 24 GeV/c	85000	2%	0.6%	1.2%	~1%
J-PARC 50 GeV/c	4.1×10^5	0.9%	0.6%	1.2%	
GSI 90 GeV/c	1.2×10^6	0.6%	0.6%	1.2%	
SPS CERN 450 GeV/c	1.26×10^6	0.5%	0.6%	1.2%	

(*) Precision on $P_{br} = \varphi(\tau)$ can be increased and the error will be less than 0.6%
private communication by D. Trautmann



Expected accuracy for πK -scattering

Estimation of error sources in $\Delta/a_{1/2} - a_{3/2} / |a_{1/2} - a_{3/2}|$ based on data taken with the upgraded DIRAC setup during 12 months (20h/day)

Single-layer target

	Number of atomic pairs n_A	Relative statistical error $(a_{1/2} - a_{3/2})$	Relative theoretical error of $(a_{1/2} - a_{3/2})$ from the ratio $\tau = f(a_{1/2} - a_{3/2})$	Relative theoretical error of $(a_{1/2} - a_{3/2})$ from the ratio $P_{br} = \varphi(\tau)^{(*)}$	Error from non point-like production
PS CERN 24 GeV/c	7000	10%	1.1%	1.2%	~1%
J-PARC 50 GeV/c	1.7×10^4	7%	1.1%	1.2%	
GSI 90 GeV/c	1.4×10^5	2.5%	1.1%	1.2%	
SPS CERN 450 GeV/c	1.26×10^5	2.5%	1.1%	1.2%	

Conclusions

Present low-energy **QCD** predictions for $\pi\pi$ and πK scattering lengths

$$\pi\pi \quad \delta a_0 = 2.3\% \quad \delta a_2 = 2.3\% \quad \delta(a_0 - a_2) = 1.5\%$$

$$\pi K \quad \delta(a_{1/2} - a_{3/2}) \approx 10\%$$

Expected results of DIRAC ADDENDUM at PS CERN

$$\tau(A_{2\pi}) \rightarrow \delta(a_0 - a_2) = \pm 2\%(stat) \pm 1\%(syst) \pm 1\%(theor)$$

$$\tau(A_{\pi K}) \rightarrow \delta(a_{1/2} - a_{3/2}) = \pm 10\%(stat) \pm \dots \pm 1.5\%(theor)$$

Observation of metastable $A_{2\pi}$

DIRAC at SPS CERN

$$\tau(A_{2\pi}) \rightarrow \delta(a_0 - a_2) = \pm 0.5\%(stat) \pm 1\%(syst) \pm 1\%(theor)$$

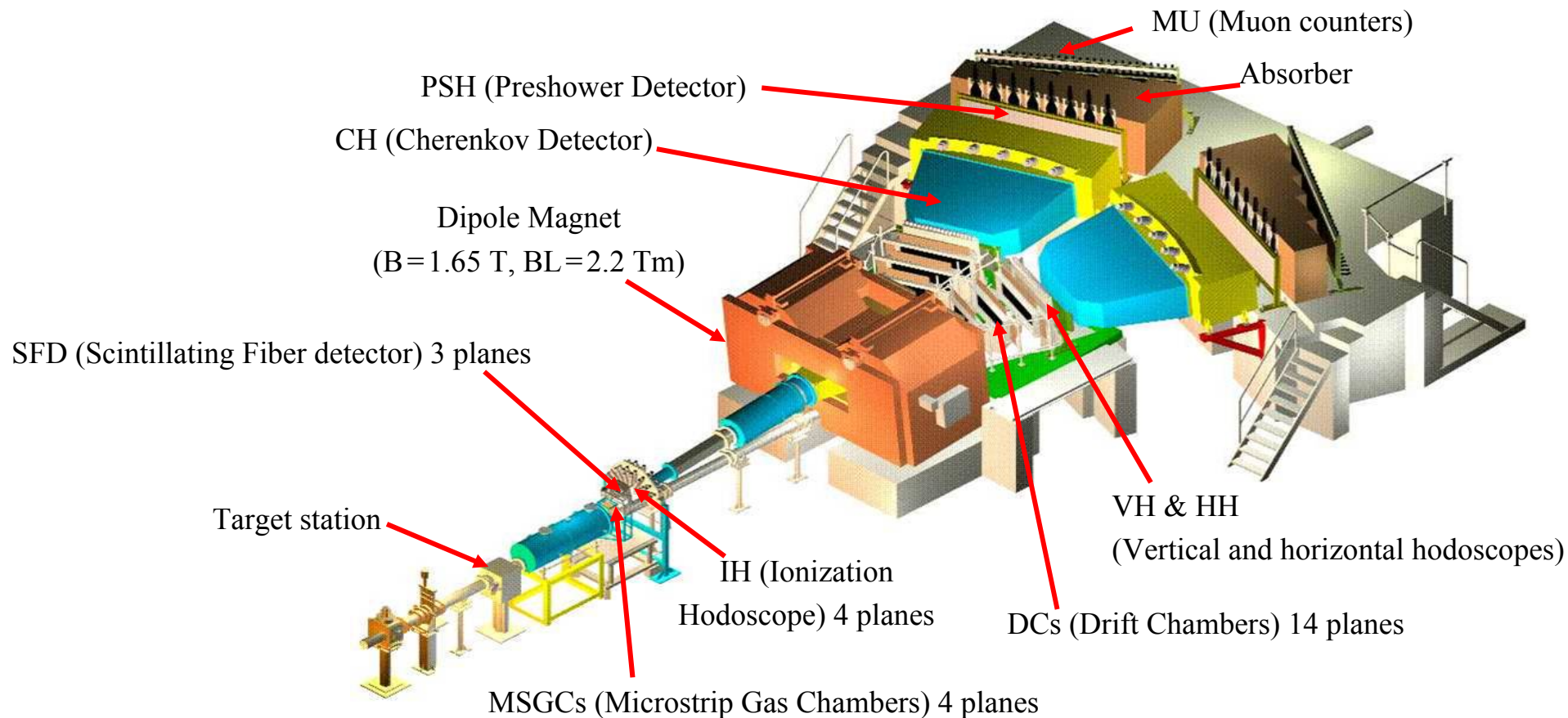
$$\tau(A_{\pi K}) \rightarrow \delta(a_{1/2} - a_{3/2}) = \pm 2.5\%(stat)$$

$$(E_{np} - E_{ns})_{\pi\pi} \rightarrow \delta(2a_0 + a_2) \approx \pm 2.5\%(stat)$$

$$(E_{np} - E_{ns})_{\pi K} \rightarrow \delta(2a_{1/2} + a_{3/2})$$

**Possibility of the observation of $(\pi^\pm \mu^\mp)$ – atoms
and of $(K^+ K^-)$ – atoms will be studied**

DIRAC isometric view



Setup features:

angle to proton beam

$$\Theta = 5.7^\circ \pm 1^\circ$$

channel aperture

$$\Omega = 1.2 \cdot 10^{-3} \text{sr}$$

momentum range

$$1.2 \leq p_\pi \leq 8 \text{ GeV}/c$$

momentum resolution

$$\Delta p/p \approx 3 \cdot 10^{-3}$$

resolution on relative momentum

$$\sigma Q_x \approx \sigma Q_y \leq 0.5 \text{ MeV}/c,$$

$$\text{and } \sigma Q_L \approx 0.5 \text{ MeV}/c$$

Upstream:

MSGC, SFD, IH

Downstream:

DC, VH, HH, Ch, PSh, Mu

Breakup probability

$$P_{br} = 0.452 \pm 0.023_{stat} \left. \begin{array}{l} +0.009 \\ -0.032 \end{array} \right\}_{syst} = 0.452^{+0.025}_{-0.039}$$

Summary of systematic uncertainties:

source

σ

CC-background

± 0.007

signal shape

± 0.002

multiple scattering angle $+5\%$
 -10%

+0.006

-0.013

K⁺K⁻ and $\bar{p}p$ pairs admixture

+0.000

-0.024

correlation function for non-point production

+0.000

-0.017

Total

+0.009

-0.032

DIRAC analysis

Improvements on systematic

CC background	no improvement	± 0.007
signal shape	no improvement	± 0.002
Multiple scattering	measured to $\pm 1\%$	+ 0.002 /-0.002
K^+K^-/pp_{bar} admixtures	to be measured*	+ 0.000 /-0.023
Finite size effects	to be measured** /improved calculations	+ 0.000 /-0.017
Total		+ 0.008 /- 0.030

* To be measured in 2006/2008 with new PID

** To be measured in 2006/2008 with new trigger for identical particles at low Q

Improvements on data quality by fine tunings

Adjustments of drift characteristics almost run-by-run
B-field adjustment and alignment tuning with Λ -mass
 \Rightarrow New preselection for all runs

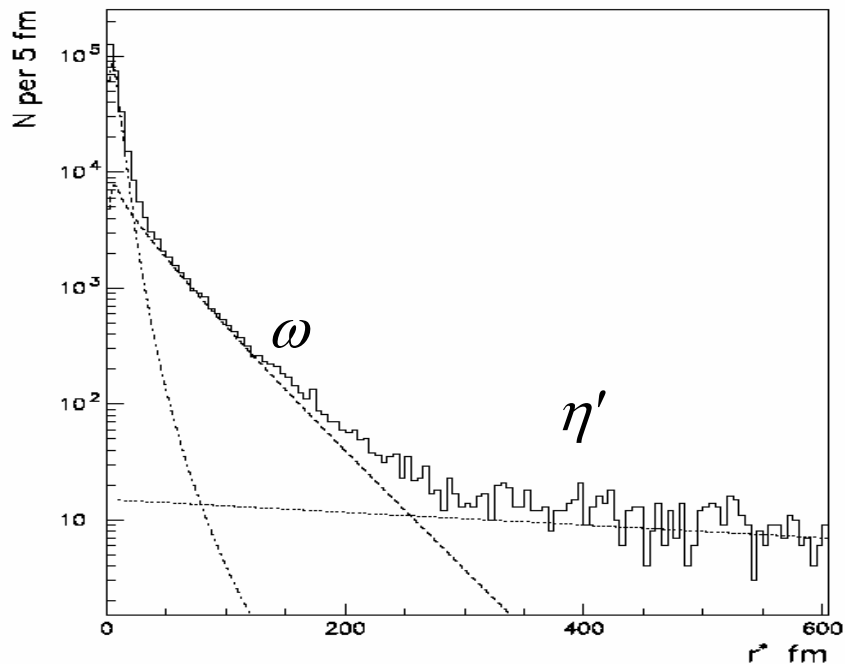
Comments on analysis strategies

Using only downstream detectors (Drift chambers) and investigating only Q_L causes less sensitivity to multiple scattering and to the signal shape. Studies are under way.



Finite-size effects (I)

- characteristic scale $|a| = 387 \text{ fm}$ (Bohr radius of $\pi\pi$ system)
- average value of $r^* \sim 10 \text{ fm}$
- range of $\omega \sim 30 \text{ fm}$
- range of $\eta' \sim 900 \text{ fm}$
- critical region of $r^* \sim |a|$ is formed by ω and η' pairs



UrQMD simulation pNi 24 GeV:

- $\sim 15\%$ ω pairs
 - $< 1\%$ η' pairs
- \Rightarrow shift in P_{br} mainly due to ω pairs

Experimental status on πK

In the 60's and 70's, set of experiments were performed to measure πK scattering amplitudes. Most of them were done studying the inelastic scattering of kaons on protons or neutrons, and later also on deuterons.

The kaon beams used in these experiments had energies ranging from 2 to 13 GeV.

The main idea of those experiments was to determine the contribution of the One Pion Exchange (OPE) mechanism. This allows to obtain the πK scattering amplitude.

Analysis of experiments gave the phases of πK -scattering in the region of $0.7 \leq m(\pi K) \leq 2.5$ GeV.

The most reliable data on the phases belong to the region $1 \leq m(\pi K) \leq 2.5$ GeV.